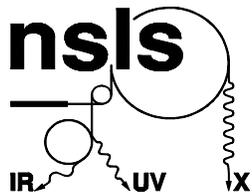


# Intense Coherent THz Pulses from the DUV-FEL Linac: Characteristics and Potential Applications

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Henrik Loos, Brian Sheehy and G. Lawrence Carr  
*National Synchrotron Light Source*  
*Brookhaven National Laboratory*

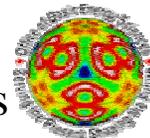
DUV-FEL Workshop, February 19th&20th, 2004



Funded under contract: DE-AC02-98CH10886



U.S. Department of Energy  
Office of Basic Energy Sciences



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NATIONAL LABORATORY

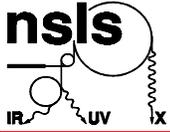
BROOKHAVEN SCIENCE ASSOCIATES

# Why THz Pulses?

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$$1 \text{ THz} = 10^{12} \text{ Hz} ; \lambda = 300 \mu\text{m}, \nu = 33 \text{ cm}^{-1}$$

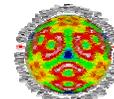
- Ultra-fast dynamics (< 1 ps time scale)
  - electronic excitations
  - magnetic excitations
- Low frequency, non-linear properties of materials
  - vibrational modes (local phonon modes)
  - electronic modes (nanoparticles / quantum dots)
- Structural transitions
  - Large E-field to coherently “shove” atoms
- Imaging
  - non-ionizing
  - crossroads between depth of penetration, chemical info.



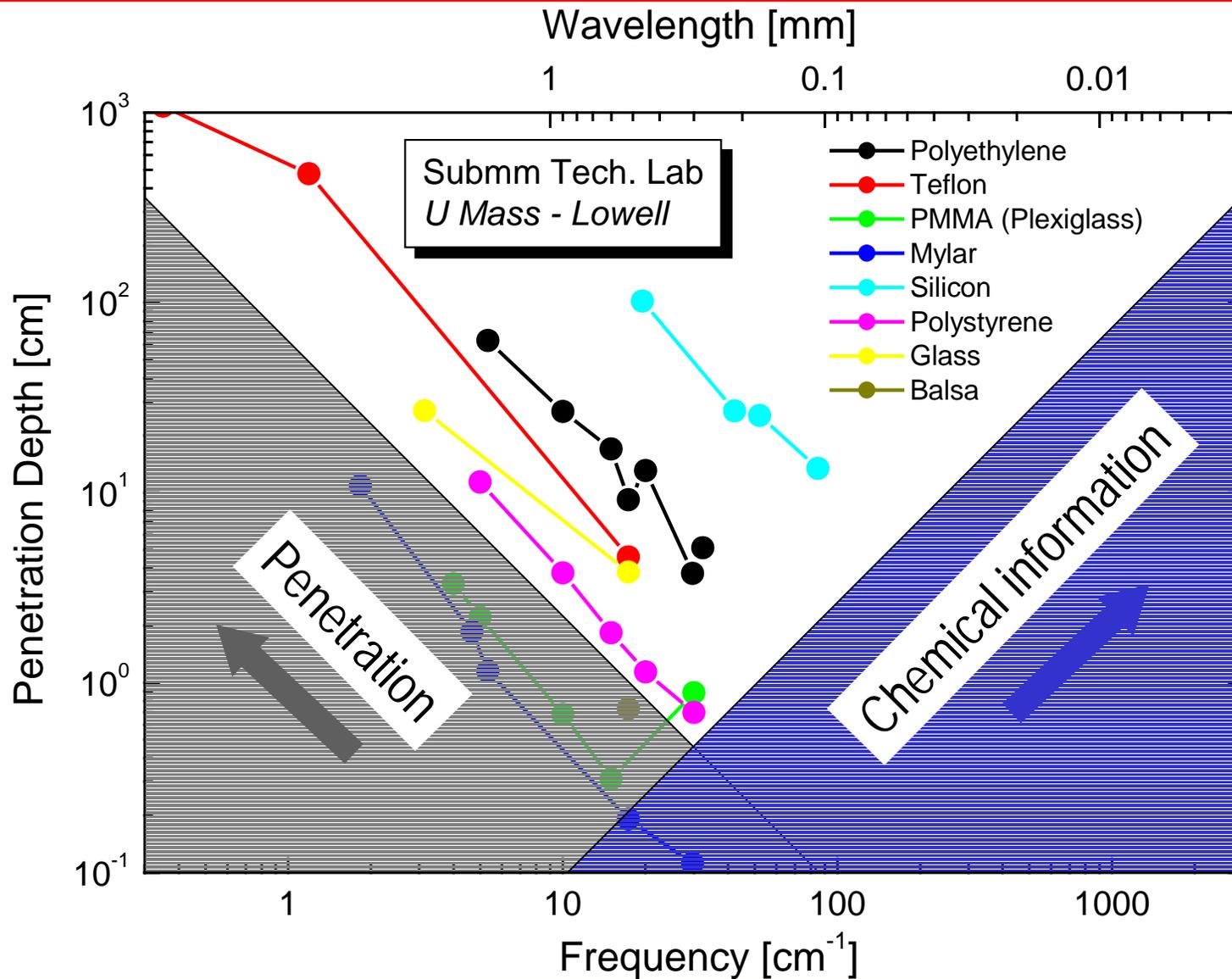
# Pulsed THz Source Requirements

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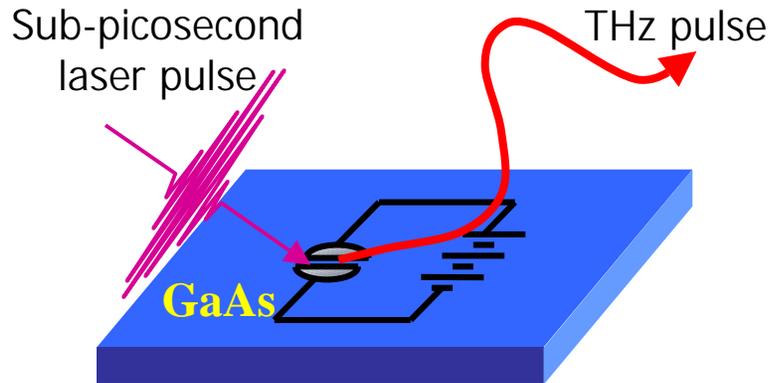
- Broad spectral content (up to at least 3 THz)
- Peak E-field approaching 1 MV/cm
- Tailored waveform
  - “Half-cycle’ pulse, other
- Two colors
  - a tunable, low energy pump pulse
  - similar or even lower energy spectroscopic probe pulse
- Sufficient average power for imaging large areas at high (video) rates.



# THz Imaging: Absorption in Common Dielectric Materials



# Energy/Power for “Table Top” THz Pulses



Radiation from acceleration of photocarriers.

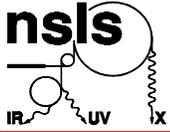
$$P = \frac{2e^2 a^2}{3c^3} \gamma^4 \quad (1 e)$$

$$10^6 \text{ V/m} \Rightarrow a \sim 10^{19} \text{ cm/s}^2$$

$$\text{Energy} = 6 \times 10^{-32} \text{ J per electron in 1 ps}$$

$$\text{or } 10^{-32} \text{ J/cm}^{-1}/\text{electron}$$

- Multiparticle ( $N^2$ ) coherent enhancement (spots size smaller than diff. limit).
- Typical energies per pulse are  $\sim 1 \text{ nJ}$  from a conventional photoconductive switch and an amplified, 250 kHz rep rate drive laser.
- Optical rectification output is similar, but extends spectra to higher frequencies.
- Higher peak power lasers (at or below 1 kHz rate) and biased antenna structures have enabled  $1 \mu\text{J}$ . (You et al, *Opt. Lett.* **18**, 290 (1993), Budiarto et al, *J. Quant. Elect.* **32**, 1839 (1996).
- Recently, Riemann et al [*Opt. Lett.* **28**, 471 (2003)] demonstrated multi-cycle mid-IR ( $\sim 800 \text{ cm}^{-1} = 24 \text{ THz}$ ) pulses with E-field of  $1 \text{ MV/cm}$ . (mixing in GaSe)



# Transition Radiation from Relativistic Electron

Transition radiation occurs when an electron crosses the boundary between two different media. For a relativistic electron ( $\beta \equiv v/c \cong 1$ ) incident on a perfect conductor, the number of photons emitted per solid angle and wavelength range is:

$$\frac{dN}{d\lambda d\Omega} = \frac{\alpha}{\pi^2 \lambda} \frac{\beta^2 \sin^2 \theta \cos^2 \theta}{(1 - \beta^2 \cos^2 \theta)}$$

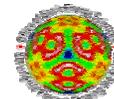
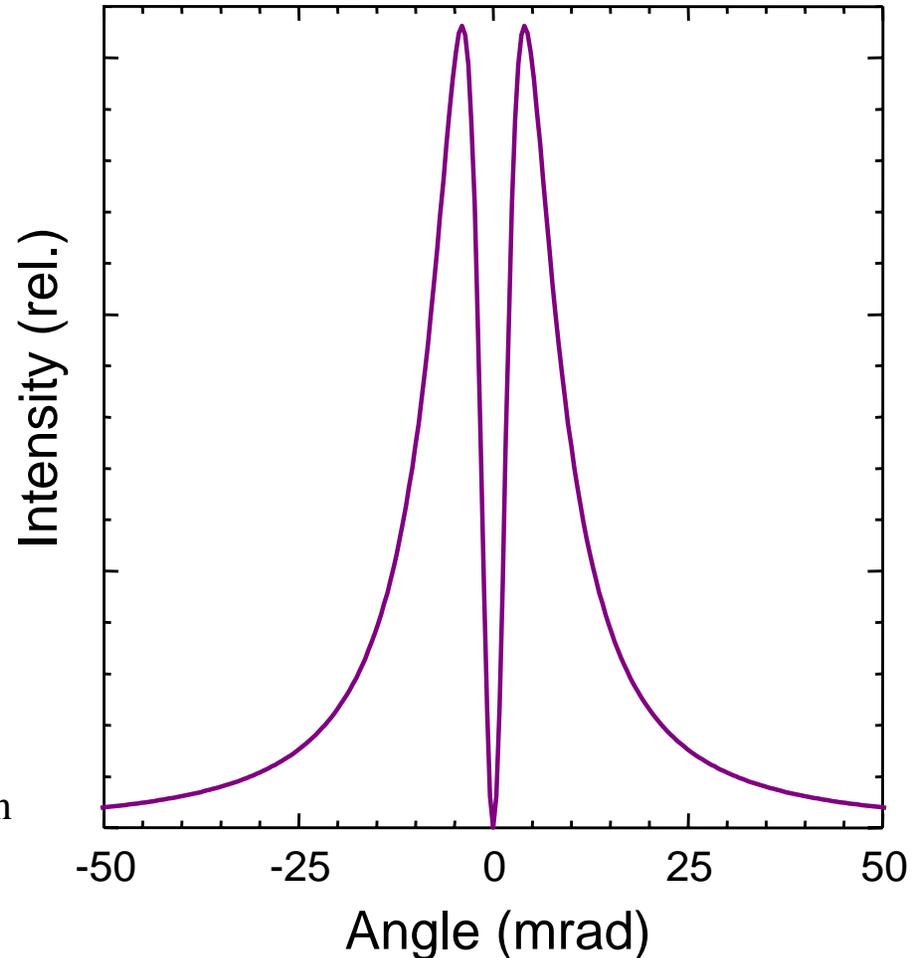
Intensity is 0 on axis, peaks at  $\theta \sim 1/\gamma$ .

Polarization is radial

$$\frac{dP}{d\bar{v}} \approx 4.61 \times 10^{-26} \left( \ln \frac{2}{1 - \beta} - 1 \right) \text{ J/cm}^{-1} \text{ per electron}$$

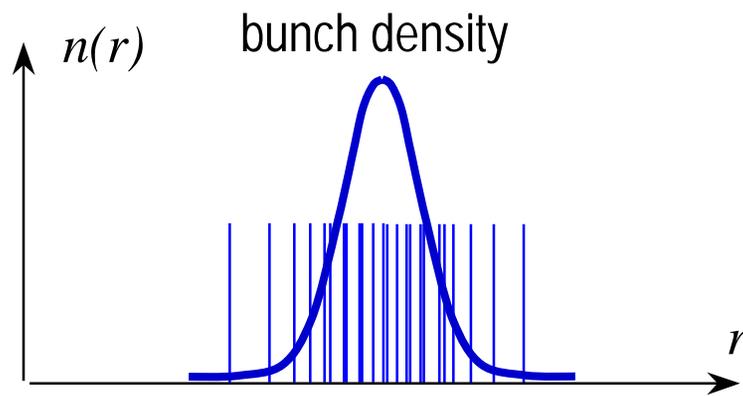
11.4 for 130 MeV  
20 for 9 GeV

Far field distribution for  $\gamma = 200$



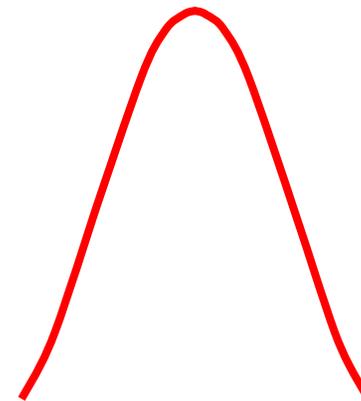
# Multi-particle Coherent Synchrotron Radiation (CSR)

- What are requirements for CSR from electrons in a bunch?
  - bunch (or some portion of it) has density variations on length scale comparable to wavelength.



$$\lambda \ll l_b$$

$$E \sim N^{1/2}; I \sim N$$



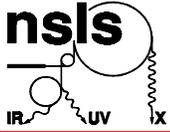
$$\lambda \gg l_b$$

$$E \sim N; I \sim N^2$$

$$\frac{dI(\omega)}{d\omega} \text{ multiparticle} = [N + N(N-1)f(\omega)] \frac{dI(\omega)}{d\omega}$$

$$\text{where } f(\omega) = \left| \int_{-\infty}^{\infty} e^{i\omega \hat{n} \cdot \vec{r}/c} S(r) dr \right|^2 \quad (\text{Nodvick \& Saxon})$$

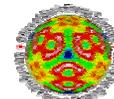
In some accelerators, bunch lengths are 100s of fs ( $\Rightarrow$  THz), and  $N$  can be large e.g.  $\sim 10^{10}$



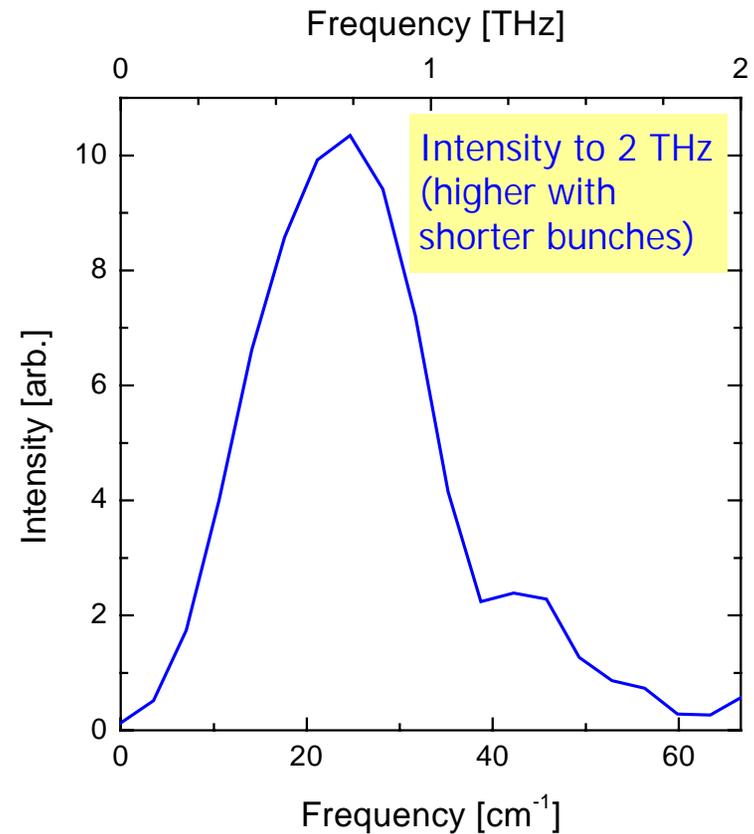
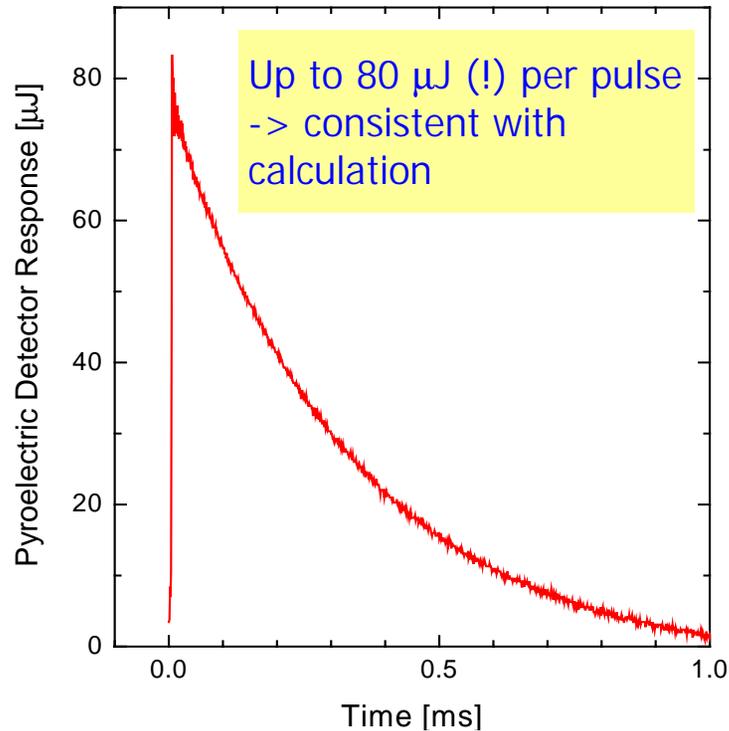
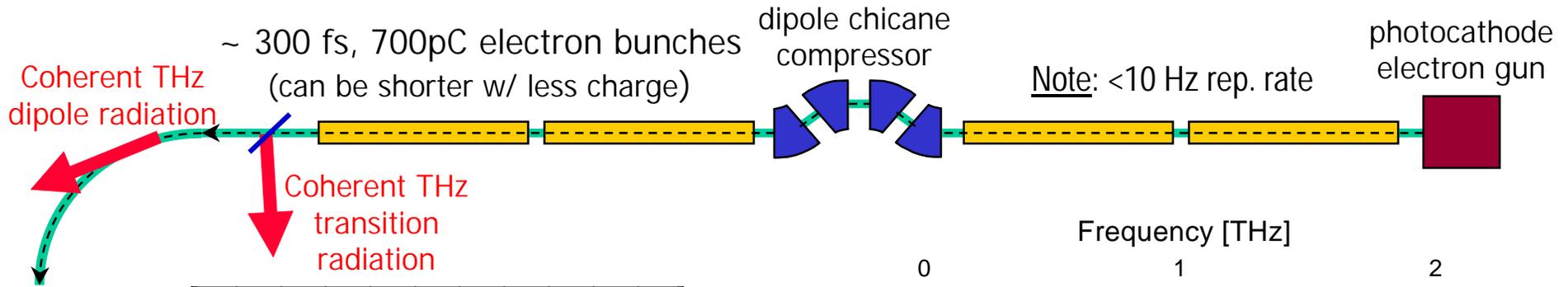
# Coherent Synchrotron Radiation (CSR)

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- 1st observations in linacs:
  - Nakazato et al (PRL '89), Happek et al (PRL '91)
- As a linac bunch diagnostic:
  - Shibata et al (PRE '94), Lai et al (PRE '94), Yan et al (PRL '00)
- As a THz source
  - Ishi et al (PRA '91), Takahashi et al (RSI '98), Carr et al., (Nature '02)
- CSR also from storage rings
  - Arpe et al, Carr et al, Anderson et al, Abo-Bakr et al. ...

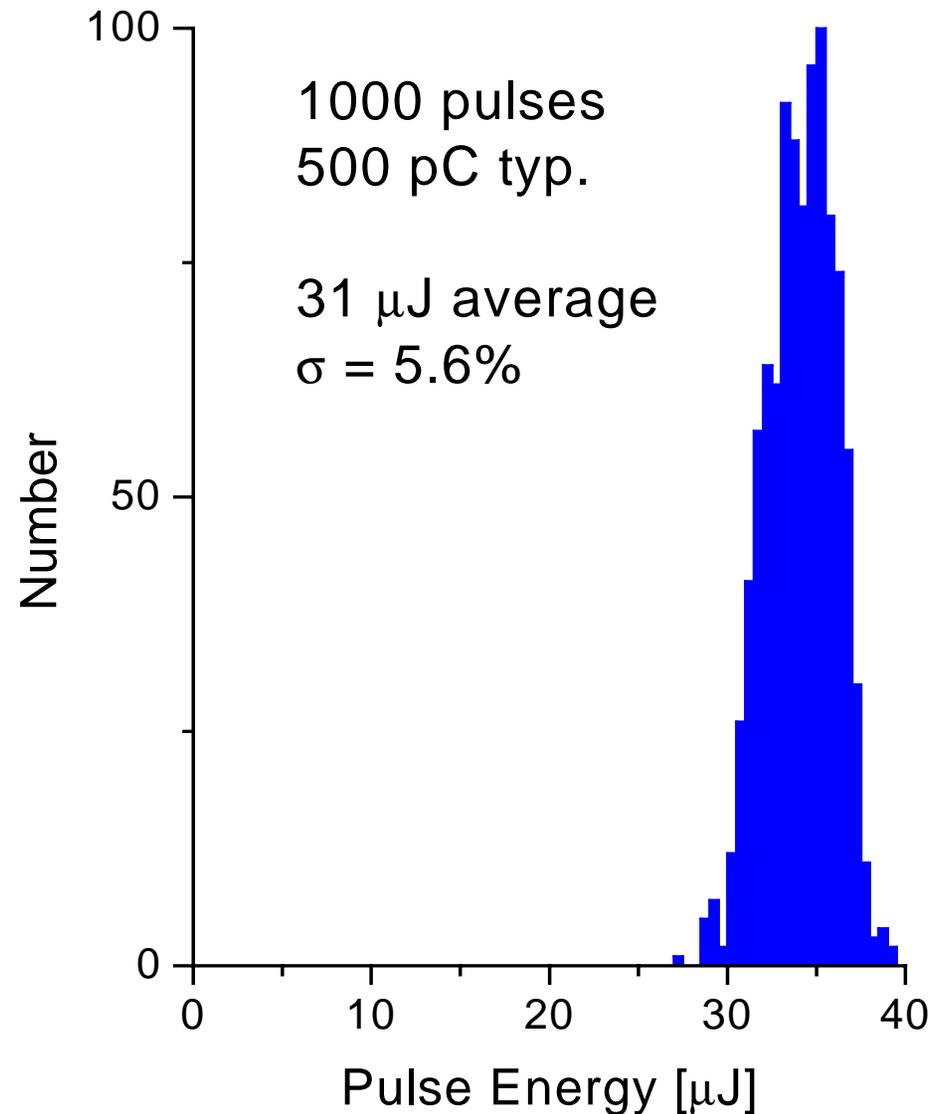


# Coherent Transition Radiation Pulses from the SDL Linac

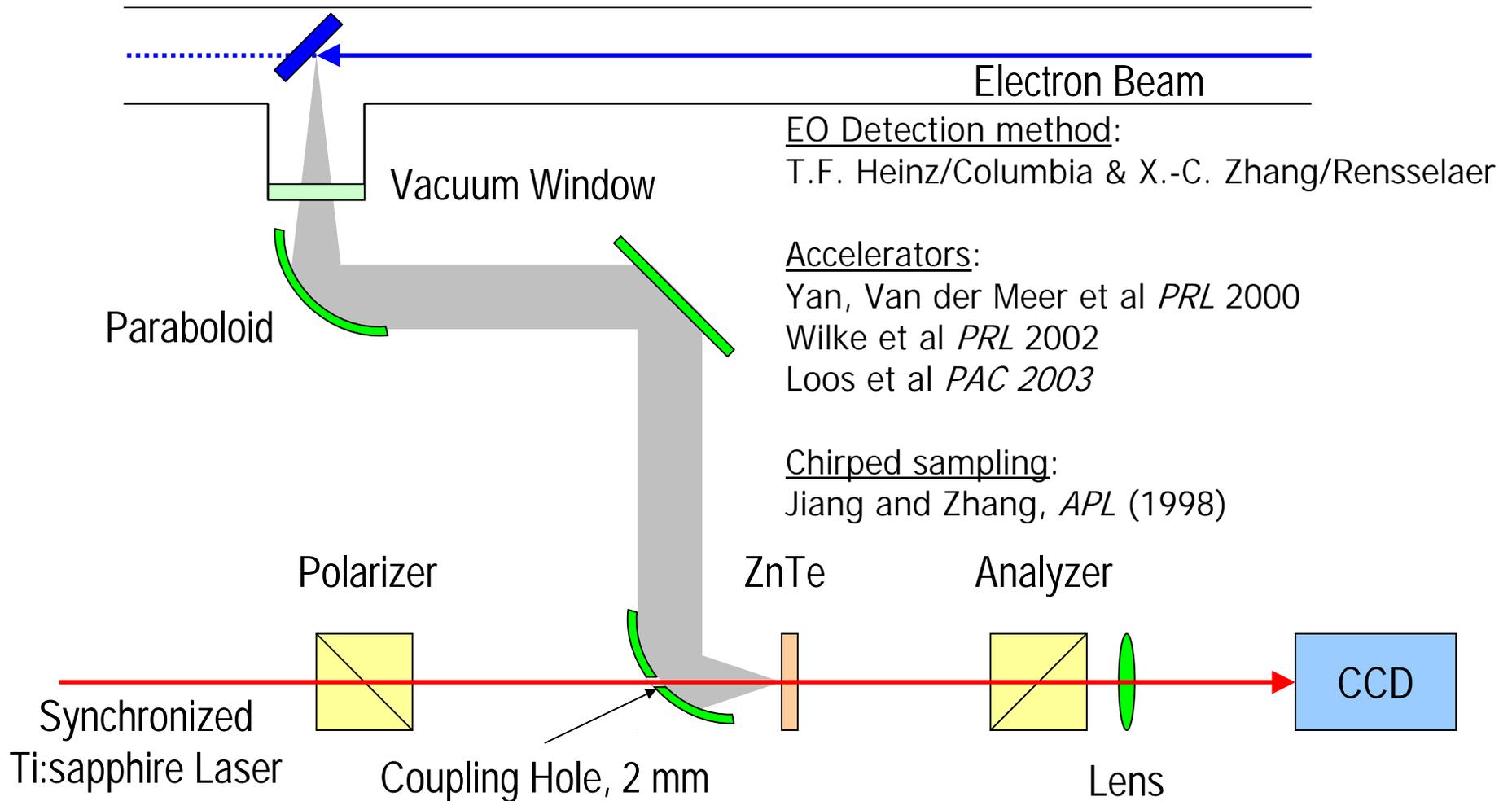


# Statistics

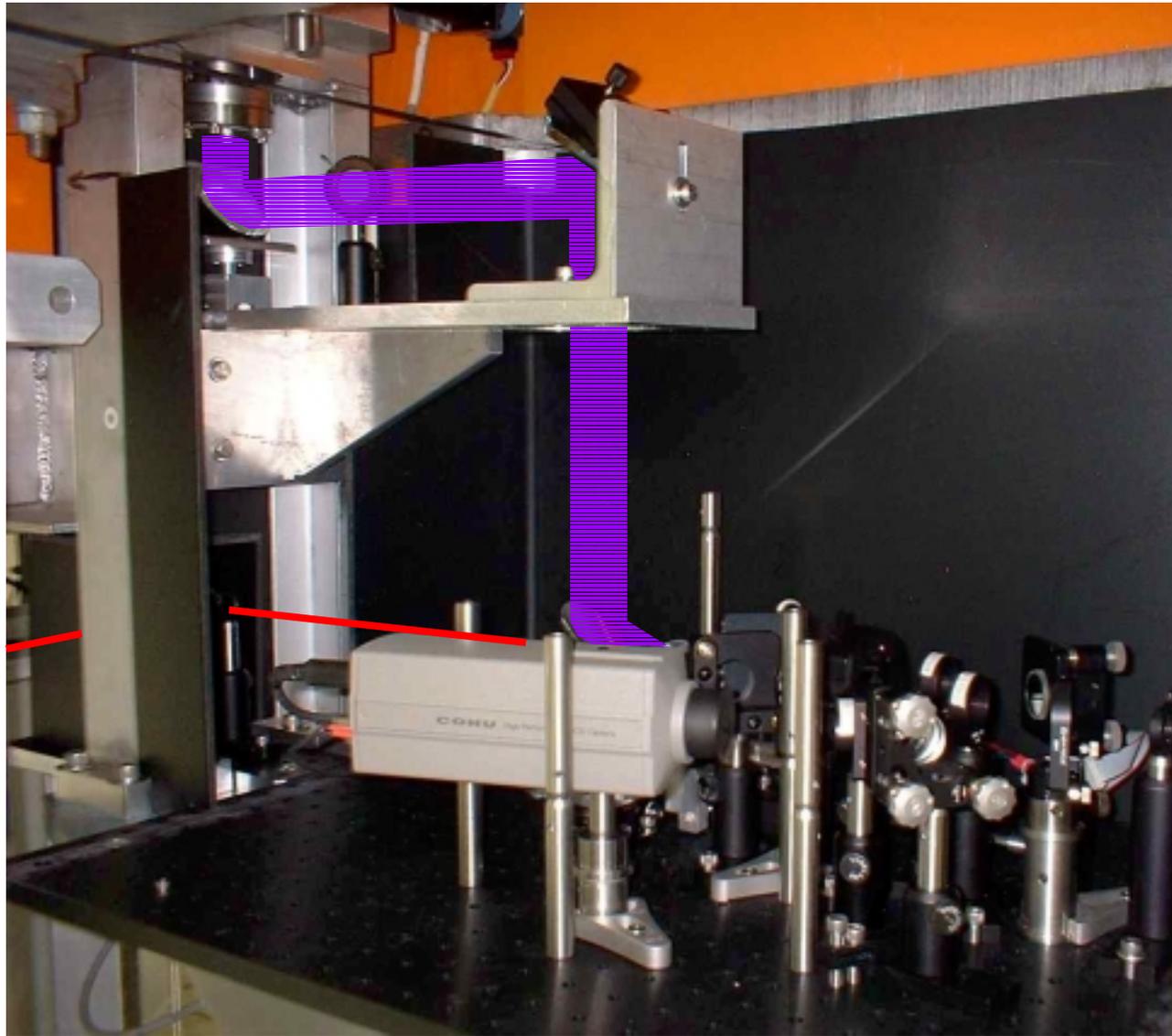
- Shot-to-shot fluctuations affect pulse usefulness for some applications.
  - Interferometry
  - Pump-probe spectroscopy
- Typical fluctuations 4 to 6% RMS.
  - due mostly to variations in charge (particle number).
- High rep. rate (average) or single-shot capability needed.



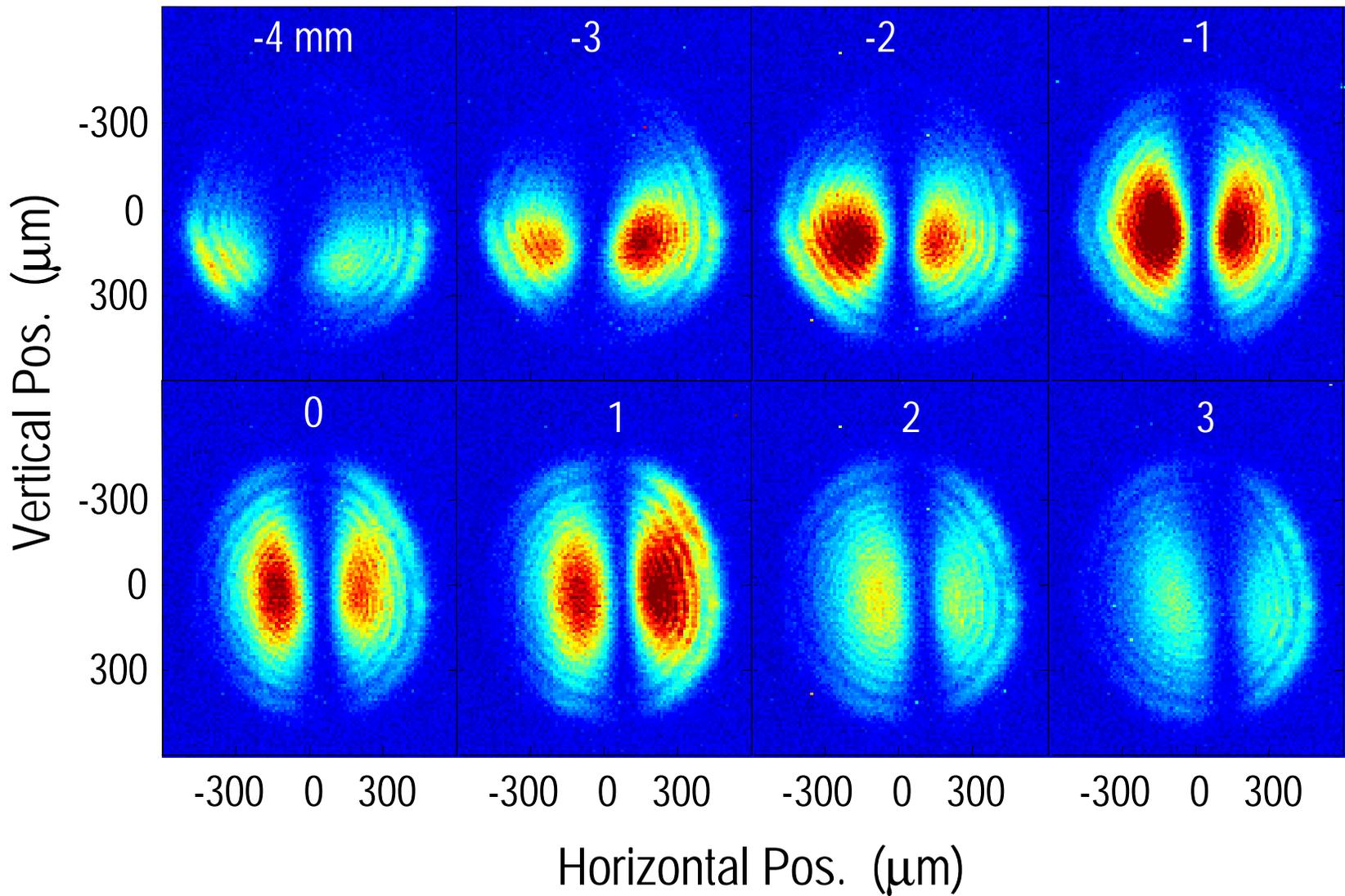
# Electro-Optic THz Pulse Detection Setup



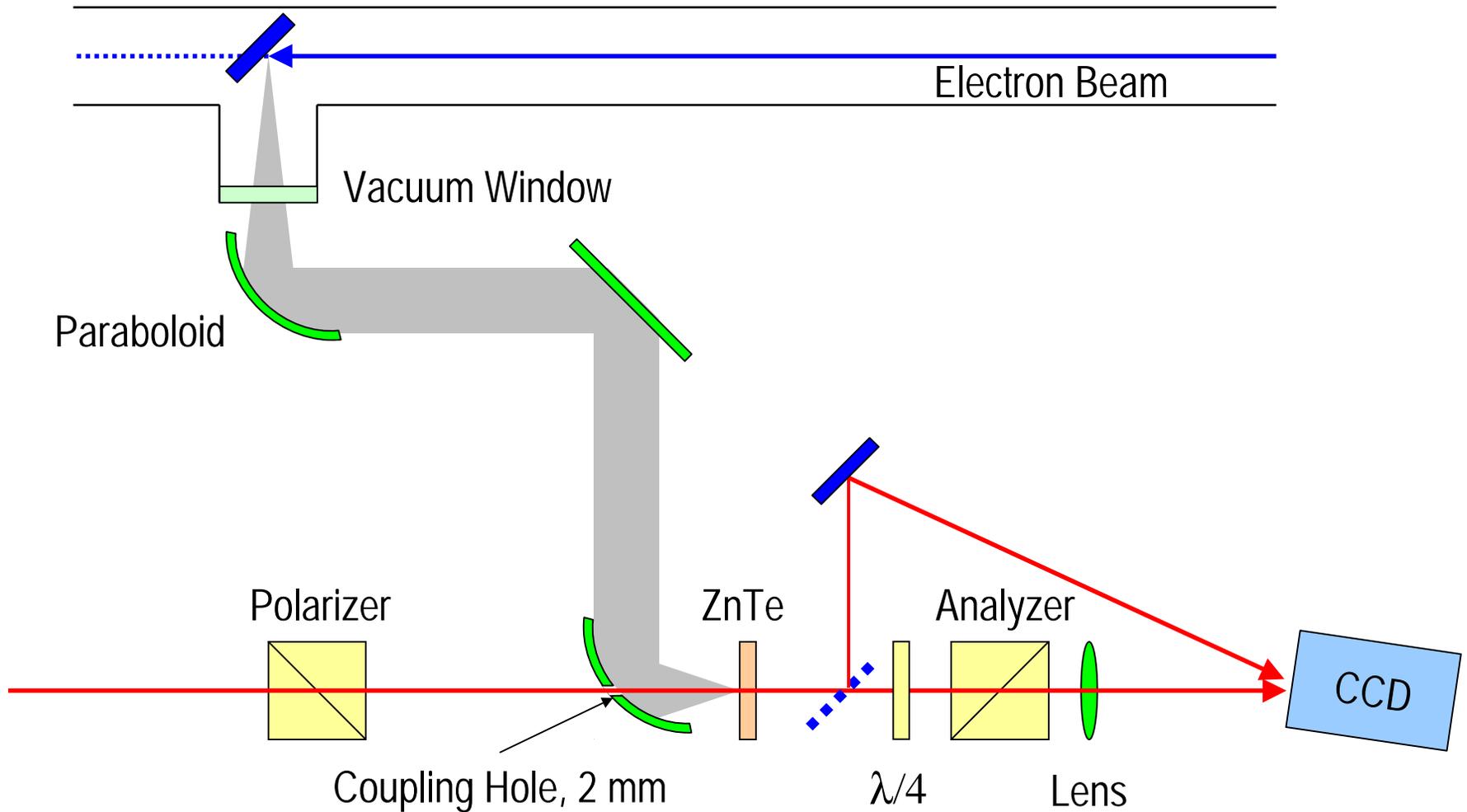
# THz and Sampling Laser Beam Path



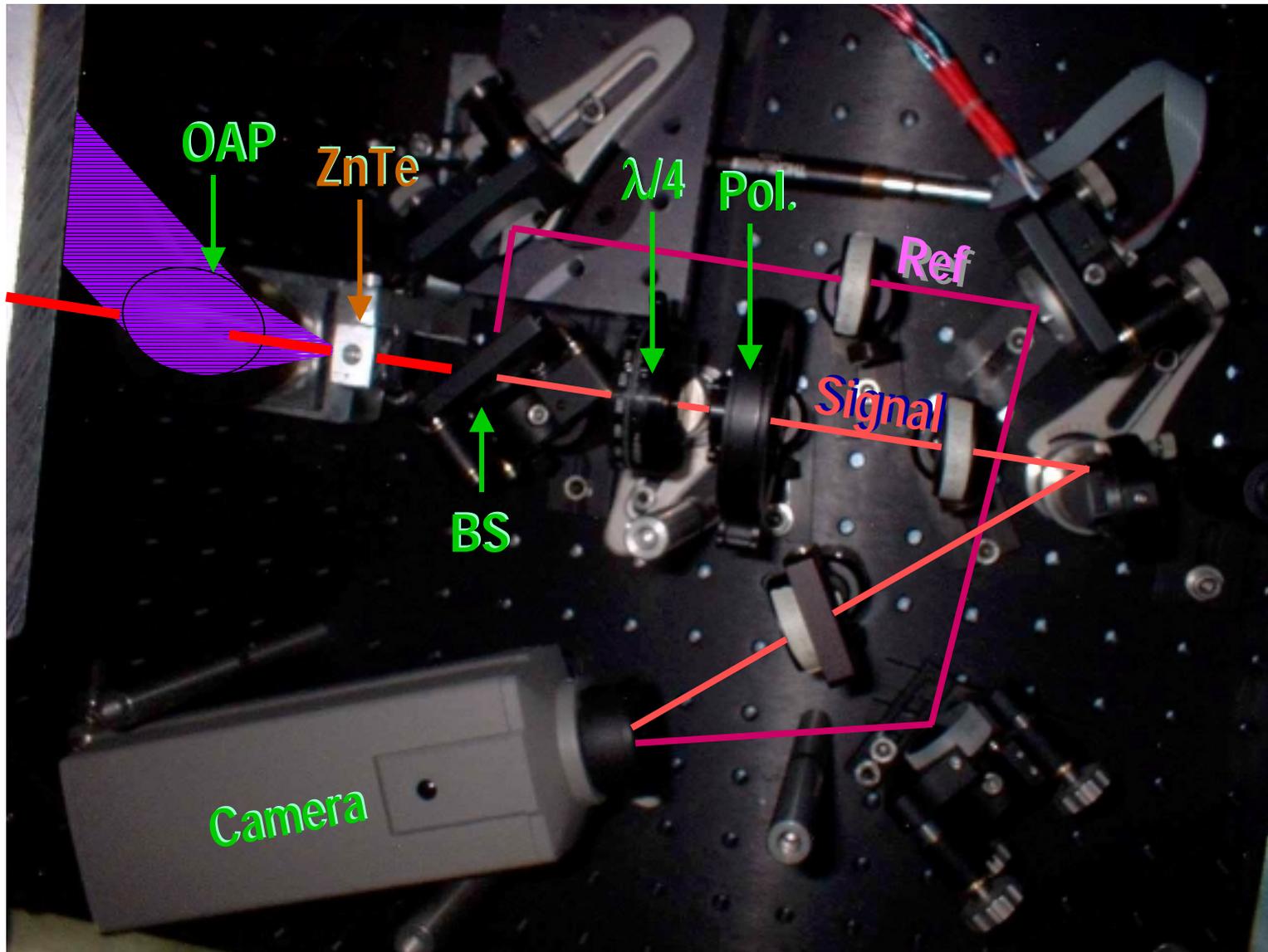
# EO Images of Electric Field (magnitude) Through Focus



# Electro-Optic THz Detection with $\lambda/4$



# Signal and Reference



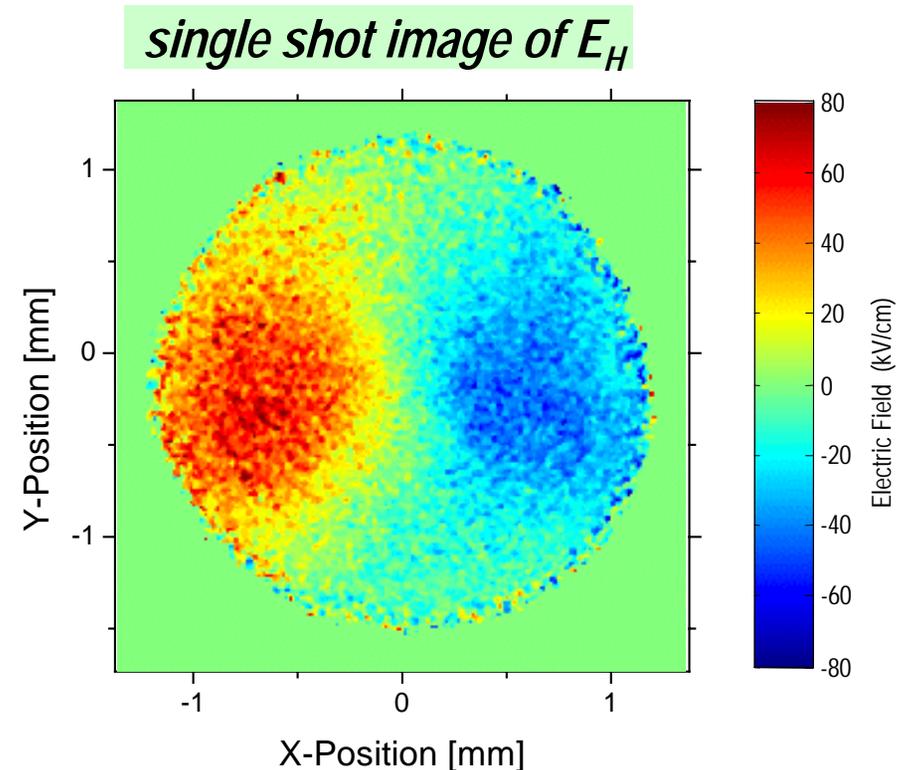
# EO Detection of SDL Linac Coherent THz Pulses

Focusing a 100  $\mu\text{J}$  pulse, 1 THz (nominal) pulse into a 1  $\text{mm}^3$  volume yields an energy density of  $10^5 \text{J/m}^3$ , so that  $E = [2D_E/\epsilon_0]^{1/2} \sim 10^8 \text{V/m}$  ( $\sim 1 \text{MV/cm}$ ).

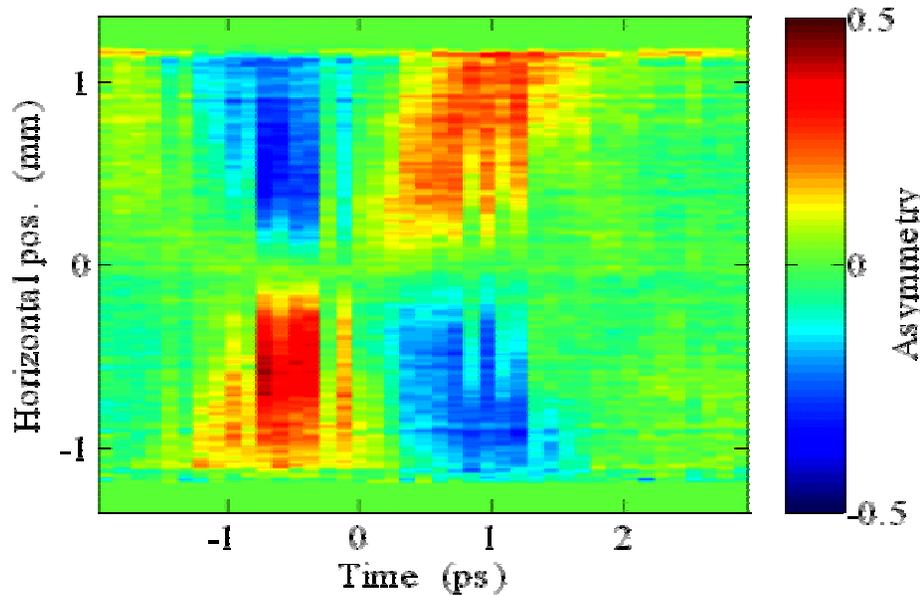
At 1 MV/cm, the magnetic field ( $B=E/c$ ) is **3 kG**.

This E-field is too large for 500  $\mu\text{m}$  ZnTe ( $E > 170 \text{kV/cm}$  yields  $> \lambda/4$  phase shift)

=> Reduce compression, lower charge to get "on-scale"



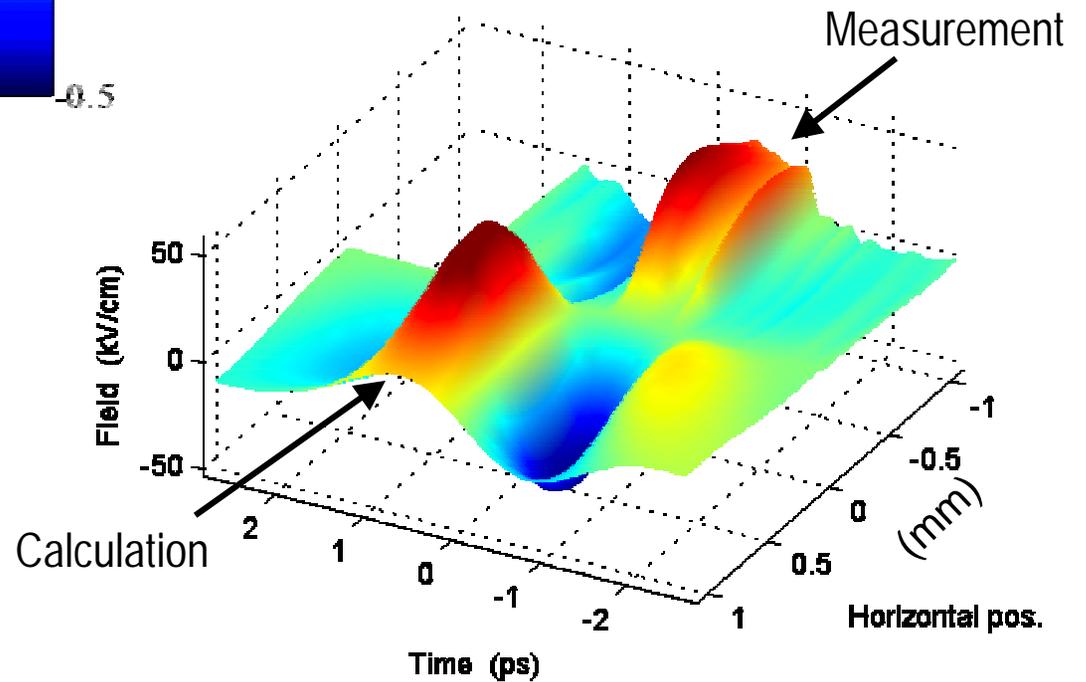
# Temporal E-Field Cross Section at Focus

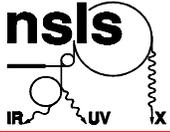


*Temporal-spatial E-field profile of coherent transition radiation pulse at  $\sim f/1.5$  focus*

Note: opposite sides are asymmetric, as shown (radial polarization)

E-field along horizontal plane

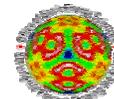




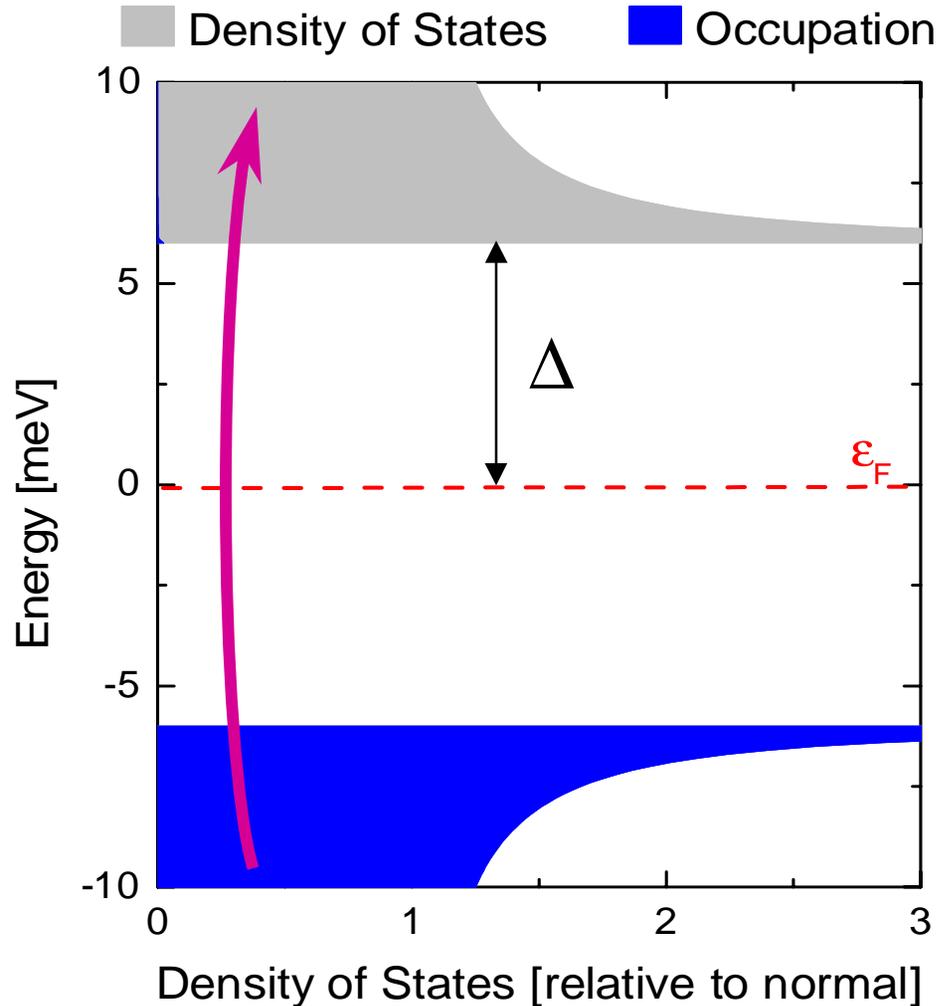
## What Can We Do with these THz Pulses?

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- ★ Use THz pulses as the pump (excitation) source in “all THz” pump probe experiments.
- ★ Study non-linear absorption in a variety of materials, including nanoparticles / quantum dots.
- ★ Induce coherent current excitations
- ★ Move atoms in their local potential wells
- ★ Orient spins in magnetic systems.



# Photoexcitation in a BCS Superconductor

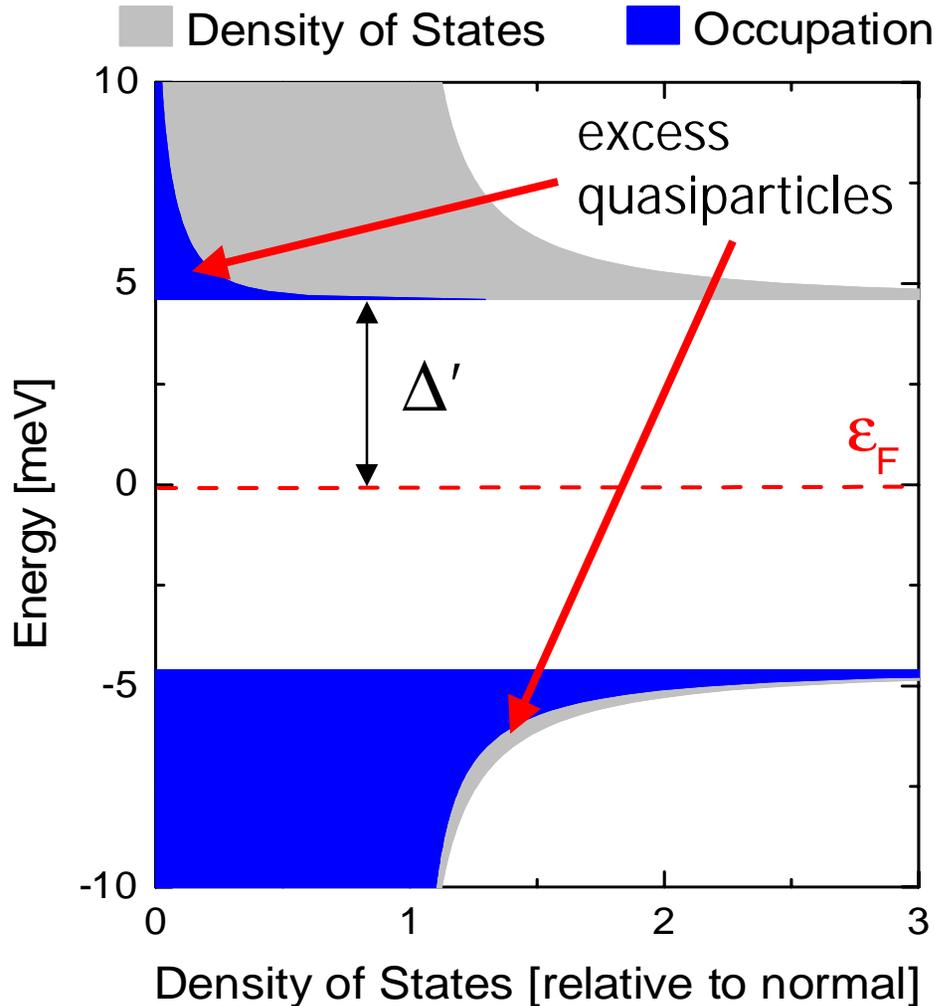


## Step 1:

Photons (a pulse from near-IR laser) break pairs, creating high energy "quasiparticle" excitations.

Photon energy 1.5 eV.

# Photoexcitation in a BCS Superconductor



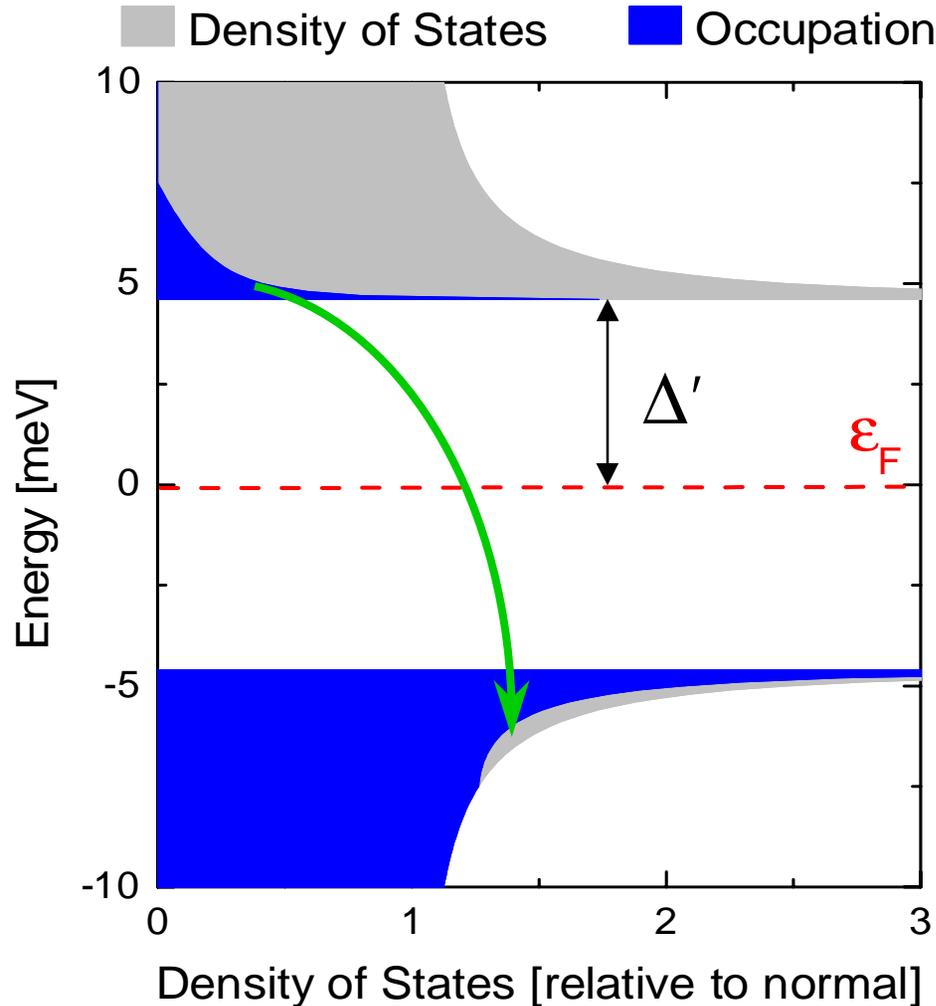
## Step 2:

High energy quasiparticles scatter and relax toward gap edge. Many more pairs broken in this process (multiplication).

Weakened superconducting state appears as reduced gap;  
 $\Delta' = \Delta - \delta$ .

Note: Quasiparticle density is out of equilibrium, but energy distribution is approximately thermalized.

# Photoexcitation in a BCS Superconductor

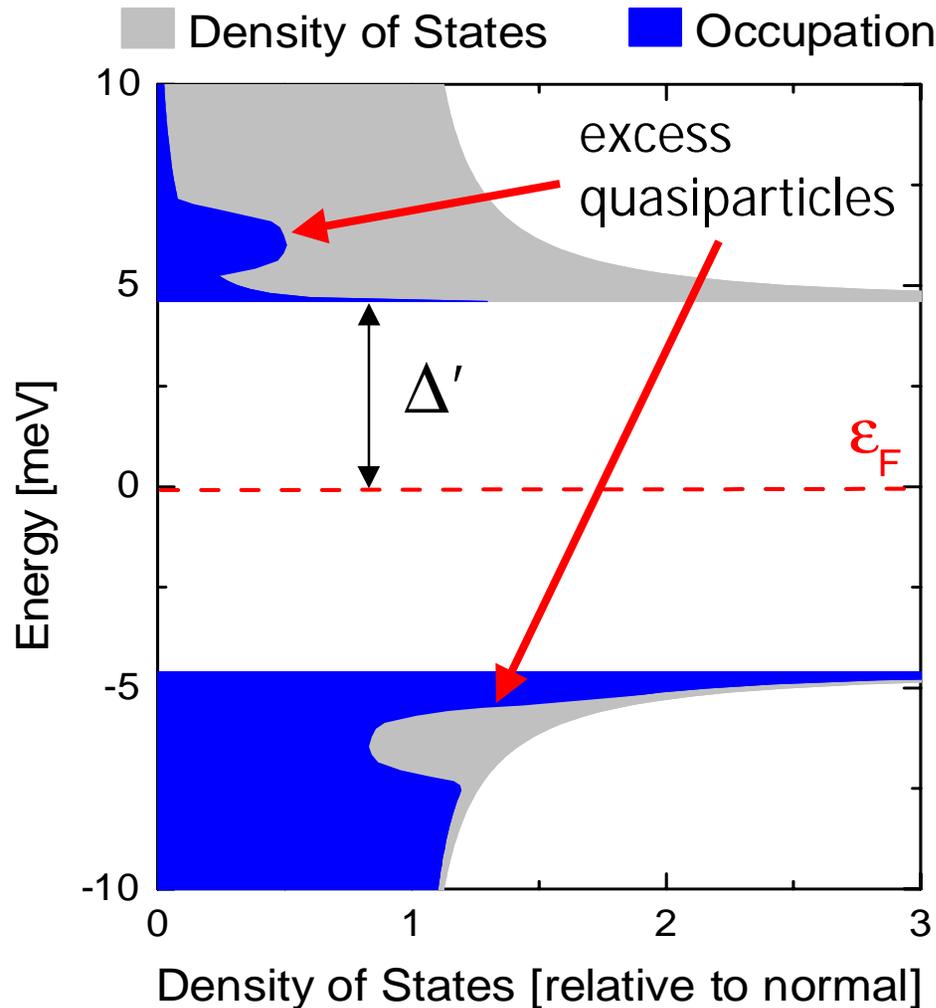


### Step 3:

Excess quasiparticles recombine to form pairs and gap is restored to full value.

Relaxation time for thin films of BCS-type superconductors is  $\sim 1$  ns.

# Photoexcitation in a BCS Superconductor



## New Experiment:

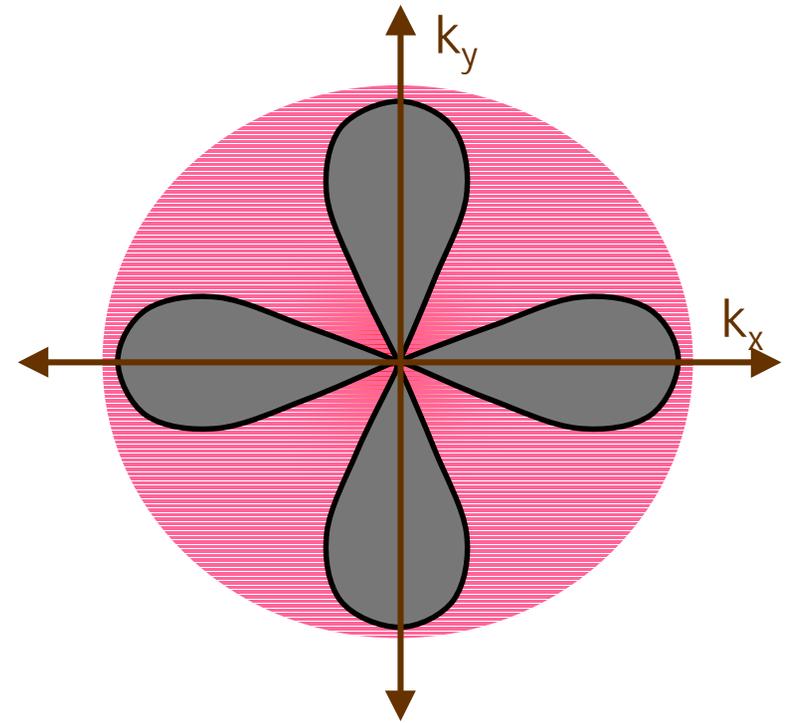
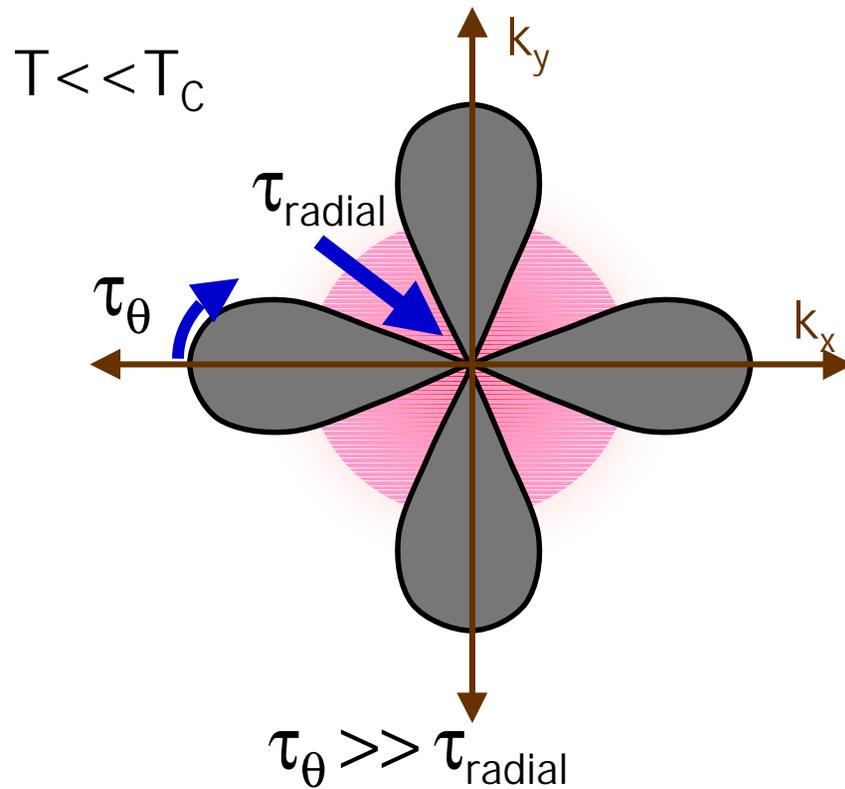
Use narrow-band, low energy pump to impose non-thermal initial distribution, probe with broadband pulse (spectroscopy).

How do these quasiparticle relax toward gap edge. How does this population affect the overall superconducting state and the energy gap?

# Pump-Probe THz Spectroscopy of a d-wave Superconductor

Near-IR laser excitation  
and fast initial relaxation

Narrow-band THz excitation  
to limit starting population energy



Goal: study  $\tau_{relax}$  as a function  
of  $E_{pump}$  (pump photon energy)

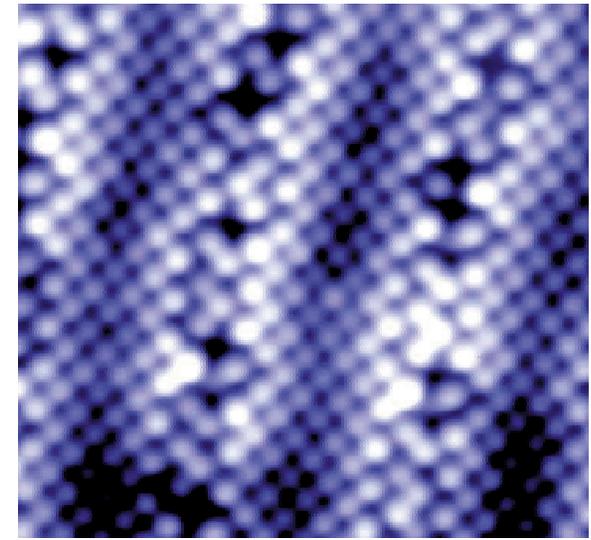
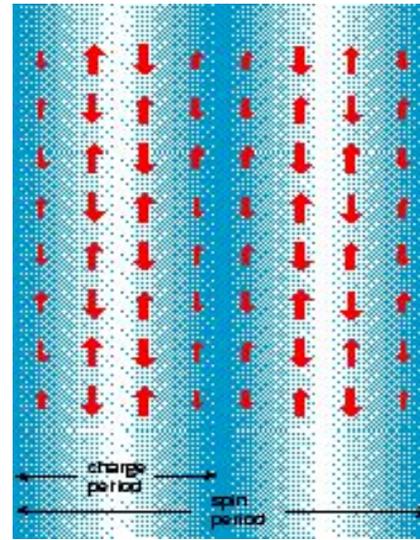
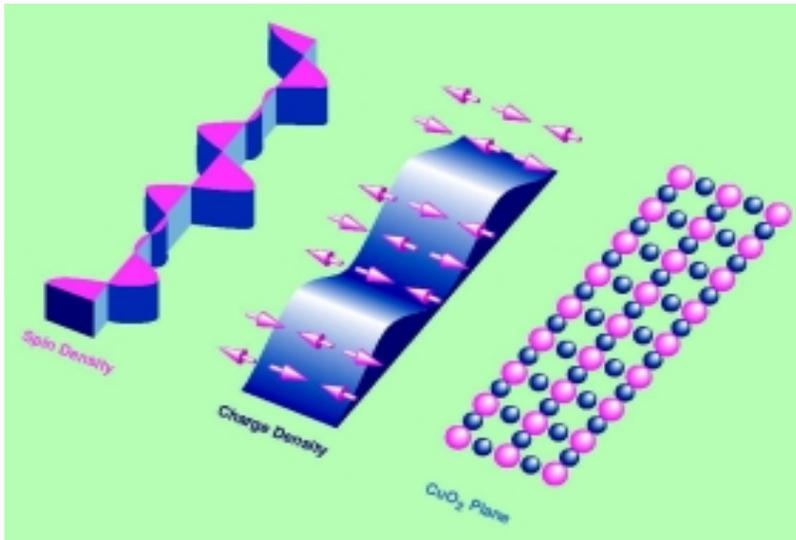
J. Orenstein et al, PRL UC Berkeley  
P.J. Hirschfeld et al, PRB U. Florida

# THz Pulse-driven Superfluid Excitations

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- Can a THz pulse induce currents that exceed the critical current in a superconductor?
  - A new type of transient excitation (avoids complicated heating effects for a DC experiment).
  - What does this superconductor look like? Is there any evidence for an energy gap or other indication of superconductivity?
  - How does it recover/relax?
- Estimates for the necessary E-field:
  - $E = J/\sigma = 10^8 \text{ A/cm}^2 / 1000 [\Omega\text{-cm}]^{-1} = 10^5 \text{ V/cm}$  or 100 kV/cm.
- Experiment: Pump with high field THz pulse (low pass filter), probe with lower intensity, broader spectral coverage pulse.

# Dynamics of Charge / Spin Stripes in Oxides



Tranquada *et al.*, Nature (1995)

J E Hoffman et al. 2002 Science 295 466

Stripes not visible in many materials: dynamic fluctuations?

Can a high-field, half-cycle THz pulse temporarily induce a preferred orientation?

Goal: THz pump, polarized THz probe to sense induced anisotropy & relaxation.

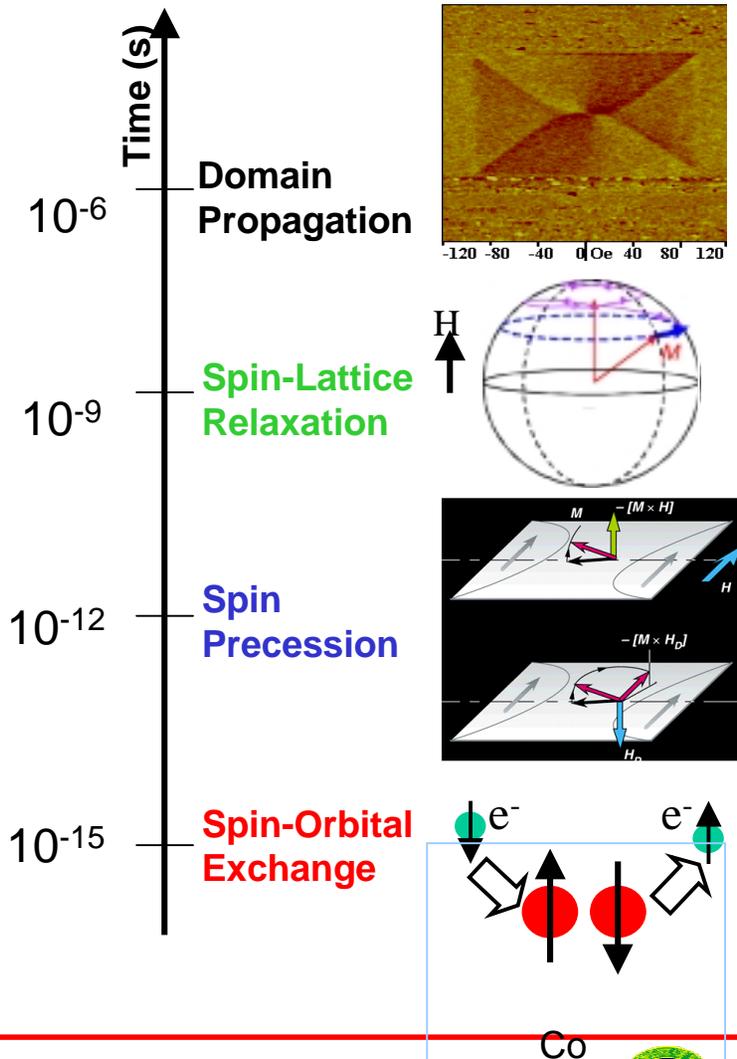
## Requirements:

Intense source of half-cycle E-field THz pulses (how big?)

THz spectroscopy probe (polarized)

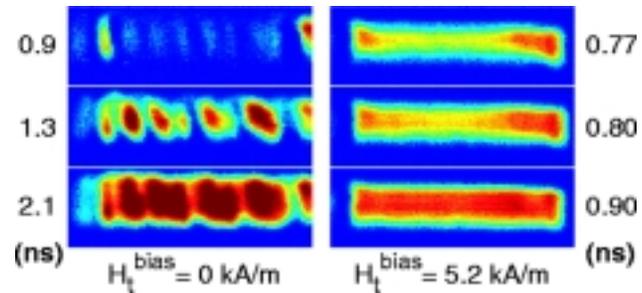
# THz Driven Magnetic Dynamics

Use ultra-short magnetic field pulses to induce spin excitations (D. Arena / NSLS)

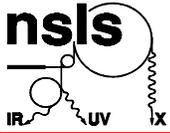


Excitation / Interaction	Timescale (sec)
Exchange interaction	10 <sup>-15</sup>
Stoner excitations	10 <sup>-15</sup> - 10 <sup>-14</sup>
Spin waves	10 <sup>-12</sup> (low q limit)
Spin - lattice relaxation	10 <sup>-12</sup> - 10 <sup>-11</sup> (in manganites)
Precessional motion	10 <sup>-10</sup> - 10 <sup>-9</sup>
Spin injection	TBD
Spin diffusion	TBD
Spin coherence	TBD

Soft Ferromagnet Dynamics Time-resolved MOKE on permalloy strip. B.C. Choi *et al.*, PRL 86, 728, (2001)



Other systems of interest: Dilute Mag. Semiconductors, Manganites.



# Opportunities in Magnetism with THz Pulses

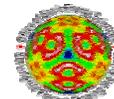
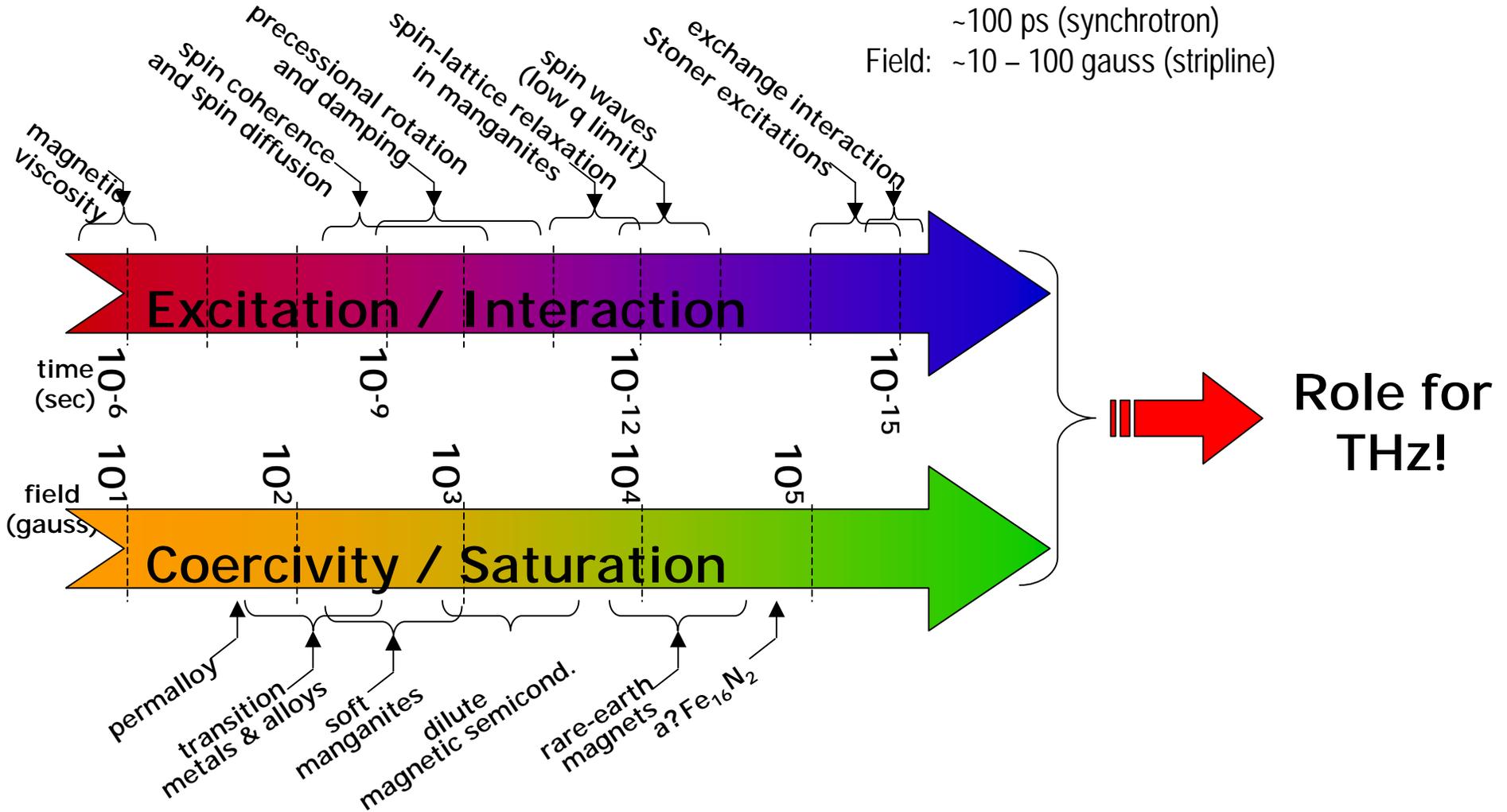
## Ultra-Short Pulses and/or High Fields -- D. Arena / NSLS

Current state of the art for "ultra-fast" dynamics experiments:

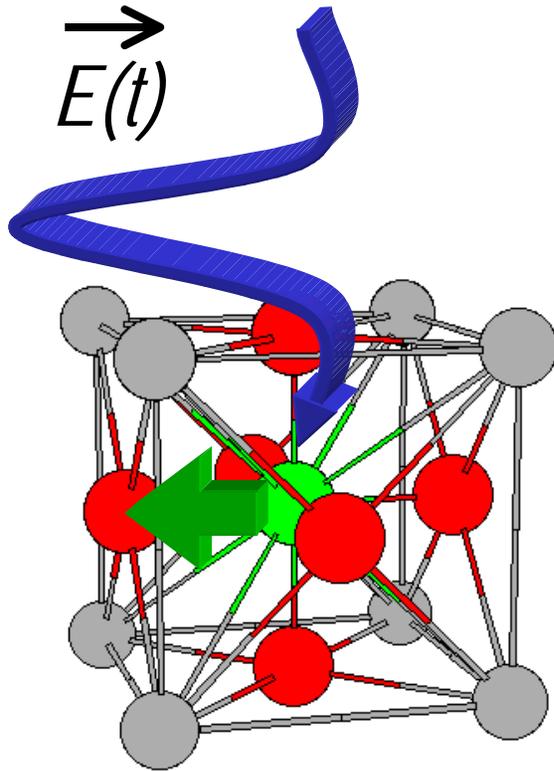
Time: ~100 fs (lasers)

~100 ps (synchrotron)

Field: ~10 – 100 gauss (stripline)



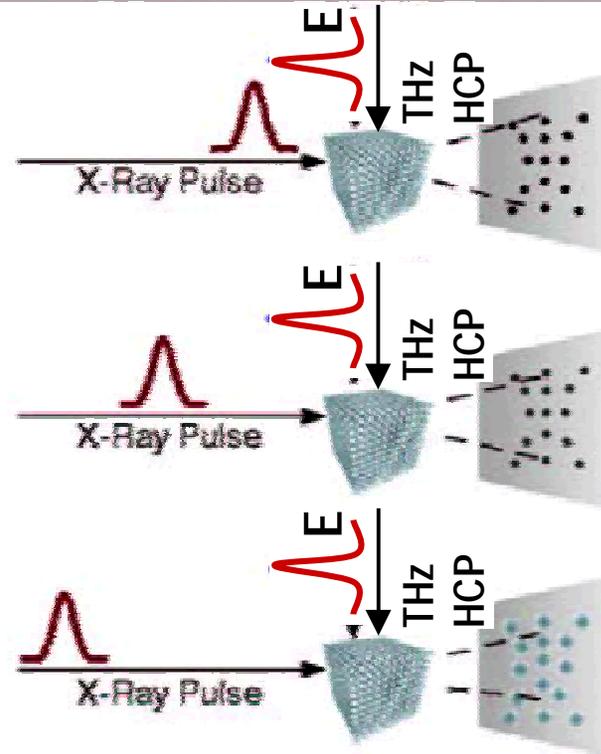
# Soft Modes in Ferroelectrics & Perovskites ( $\text{PbTiO}_3$ )



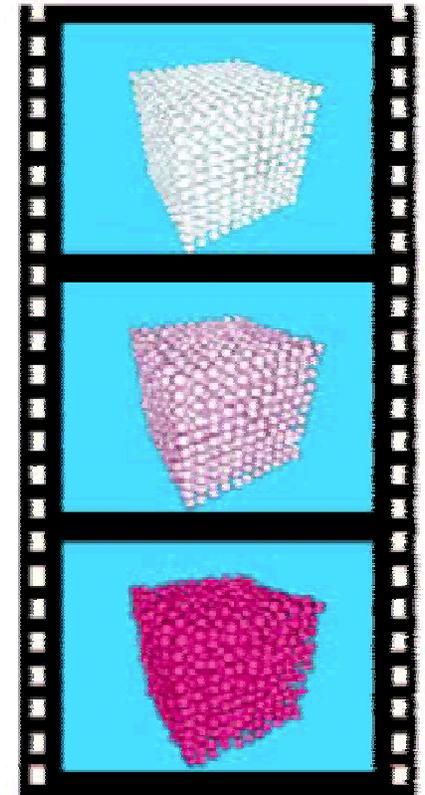
Use half-cycle pulse to coherently drive atoms, probe motion as a function of time (needs diffraction probe).

DOE Workshop on Ultrafast X-rays

## Ultrafast X-Ray Diffraction



## Movie of Atomic Movement



Observe "shift" in diffraction spot(s)

# Summary

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- Per particle, relativistic electrons offer a significant intensity/energy advantage over conventional pulsed THz sources.
  - *80  $\mu\text{J}$  per pulse, E-field approaching 1 MV/cm demonstrated at DUV-FEL/SDL.*
  - *High pulse energy should be sufficient to create novel coherent excitations in solids.*
  - *Radial polarization of transition radiation may be useful for coupling into coaxial guides.*
- Characteristics compatible with EO detection methods.
  - *Necessary for coherent imaging technique*
  - *Single-shot spectroscopy (chirped sampling pulse) desirable.*
- Spectral content determined by electron beam density.
  - *Can we make two colors? What shapes can be achieved?*
- Existing accelerators producing coherent THz pulses were not designed for this purpose.
  - *Opportunity to optimize parameters for greatest flexibility versus cost/size.*
  - *Shortest achievable bunch length versus charge, versus energy?*